

# Solutions to Extra Credit Problems

Charles Martin

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1. Evaluate

$$\int_0^1 \left\{ \frac{1}{x} \right\}^2 dx,$$

where  $\{\alpha\} = \alpha - \lfloor \alpha \rfloor$  denotes the fractional part of  $\alpha$ .

Answer:  $\ln(2\pi) - \gamma - 1$ , where  $\gamma$  is Euler's constant.

*Proof.* We remark that the integral exists because  $0 \leq \{1/x\}^2 \leq 1$  for all  $x \in (0, 1]$ . Fix a positive integer  $N$  and consider

$$\int_{1/N}^1 \left\{ \frac{1}{x} \right\}^2 dx.$$

Substituting  $x \mapsto 1/x$  gives

$$\int_{1/N}^1 \left\{ \frac{1}{x} \right\}^2 dx = \int_1^N \frac{\{x\}^2}{x^2} dx.$$

For each positive integer  $k$ , if  $k \leq x < k+1$  then  $\{x\} = x - k$ . Therefore,

$$\int_1^N \frac{\{x\}^2}{x^2} dx = \sum_{k=1}^{N-1} \int_k^{k+1} \frac{(x-k)^2}{x^2} dx = \sum_{k=1}^{N-1} \left( 2 - \frac{1}{k+1} - 2k \ln(k+1) + 2k \ln k \right).$$

Note that

$$\sum_{k=1}^{N-1} 2 = 2N - 2 \quad \text{and} \quad \sum_{k=1}^{N-1} \frac{-1}{k+1} = -\ln N - \gamma_N + 1,$$

where  $\gamma_N \rightarrow \gamma$  as  $N \rightarrow \infty$ . Furthermore,

$$\sum_{k=1}^{N-1} -2k \ln(k+1) + 2k \ln k = -2 \sum_{k=1}^{N-1} [(k+1) \ln(k+1) - k \ln k] + 2 \sum_{k=1}^{N-1} \ln(k+1) = -2N \ln N + 2 \ln(N!).$$

From Stirling's approximation

$$N! = \sqrt{2\pi N} \left( \frac{N}{e} \right)^N e^{\theta_N}$$

(where  $\theta_N \rightarrow 0$  as  $N \rightarrow \infty$ ), we have

$$2 \ln(N!) = 2N \ln N - 2N + \ln(2\pi N) + 2\theta_N.$$

Altogether,

$$\sum_{k=1}^{N-1} \left( 2 - \frac{1}{k+1} - 2k \ln(k+1) + 2k \ln k \right) = \ln(2\pi) - \gamma_N - 1 + 2\theta_N.$$

This gives

$$\int_{1/N}^1 \left\{ \frac{1}{x} \right\}^2 dx = \ln(2\pi) - \gamma_N - 1 + 2\theta_N,$$

and taking  $N \rightarrow \infty$  gives  $\ln(2\pi) - \gamma - 1$ . □

2. Evaluate  $\sum_{k=2}^{\infty} \frac{\zeta(k) - 1}{k}$ .

Answer:  $1 - \gamma$ , where  $\gamma$  is Euler's constant.

*Proof.* Call the sum  $S$ . Insert the series definition for  $\zeta$  and note that  $\sum_{n=1}^{\infty} n^{-k} - 1 = \sum_{n=2}^{\infty} n^{-k}$  to get

$$S = \sum_{k=2}^{\infty} \sum_{n=2}^{\infty} \frac{1}{kn^k}.$$

These terms are all positive, so an application of Tonelli's theorem (equally applicable to both sums and integrals) gives

$$S = \sum_{n=2}^{\infty} \sum_{k=2}^{\infty} \frac{1}{kn^k}.$$

Recognize the series for the logarithm:

$$-\ln\left(1 - \frac{1}{n}\right) = \sum_{k=1}^{\infty} \frac{1}{kn^k}.$$

Thus

$$S = \sum_{n=2}^{\infty} \left[ -\ln\left(1 - \frac{1}{n}\right) - \frac{1}{n} \right] = \sum_{n=2}^{\infty} \left[ \ln n - \ln(n-1) - \frac{1}{n} \right]$$

Take a large positive integer  $N$ ; we can telescope to get

$$\sum_{n=2}^N \left[ \ln n - \ln(n-1) - \frac{1}{n} \right] = \ln N - H_N + 1,$$

where we use the notation  $H_N = 1 + 1/2 + 1/3 + \dots + 1/N$ . It is well known that  $H_N = \ln N + \gamma + O(1/N)$ , so

$$S = \lim_{N \rightarrow \infty} (\ln N - H_N + 1) = \lim_{N \rightarrow \infty} (1 - \gamma + O(1/N)) = 1 - \gamma,$$

as claimed. □

3. Evaluate  $\sum_{n=0}^{\infty} \binom{2n}{n}^{-1}$ .

Answer:  $4/3 + 2\pi\sqrt{3}/27$

*Proof.* There are many approaches—using arcsine, using hypergeometric functions, etc.—but we'll use the gamma and beta functions. Define the functions

$$\Gamma(n) = \int_0^{\infty} t^{n-1} e^{-t} dt \quad \text{and} \quad B(x, y) = \int_0^1 t^{x-1} (1-t)^{y-1} dt.$$

It is easy to check (via integration by parts and mathematical induction) that  $\Gamma(x+1) = x\Gamma(x)$  in general and that  $\Gamma(n) = (n-1)!$  for integers  $n$ . Furthermore,

$$B(x, y) = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)},$$

which follows from simple manipulation of the integrals. Altogether, we have

$$\binom{2n}{n}^{-1} = \frac{(n!)^2}{(2n)!} = \frac{\Gamma(n+1)^2}{\Gamma(2n+1)} = \frac{(2n+1)\Gamma(n+1)^2}{\Gamma(2n+2)} = (2n+1)B(n+1, n+1).$$

Written in terms of the integrals,

$$\sum_{n=0}^{\infty} \binom{2n}{n}^{-1} = \sum_{n=0}^{\infty} (2n+1) \int_0^1 t^n (1-t)^n dt = \sum_{n=0}^{\infty} \int_0^1 (2n+1)(t-t^2)^n dt.$$

For  $t \in [0, 1]$  and  $n \geq 0$ , we have  $(2n+1)(t-t^2)^n \geq 0$ , so the monotone convergence theorem gives

$$\sum_{n=0}^{\infty} \binom{2n}{n}^{-1} = \int_0^1 \sum_{n=0}^{\infty} (2n+1)(t-t^2)^n dt.$$

The sum can be evaluated by differentiating the geometric series:

$$\sum_{n=0}^{\infty} x^n = \frac{1}{1-x} \quad \Rightarrow \quad \sum_{n=0}^{\infty} nx^n = \frac{x}{(1-x)^2},$$

so we have

$$\sum_{n=0}^{\infty} \binom{2n}{n}^{-1} = \int_0^1 \frac{1+t-t^2}{(1-t+t^2)^2} dt.$$

This integral can be evaluated to get the answer of  $4/3 + 2\pi\sqrt{3}/27$ . □

4. Let  $R$  be a nice region in  $\mathbb{R}^2$  containing the origin and  $p$  be the Green function of  $R$  with singularity at zero. Assume that  $\partial R$  moves with normal velocity given by  $V = \partial p / \partial n$  (with  $p$  changing so that it remains the Green function of  $R$ ). Given a harmonic function  $h$ , prove that

$$\frac{d}{dt} \iint_R h dA = h(0).$$

*Proof.* The following is a kinematic relation for any such moving region; its proof can be found in most books on fluid mechanics.

$$\frac{d}{dt} \iint_R h dA = \oint_{\partial R} hV ds.$$

Since  $V = \partial p / \partial n = \nabla p \cdot \mathbf{n}$ , this becomes

$$\frac{d}{dt} \iint_R h dA = \oint_{\partial R} h \nabla p \cdot d\mathbf{n},$$

a 2D flux integral. To proceed, we use Green's first identity:

$$\iint_R (h\Delta p - p\Delta h) dA = \oint_{\partial R} (h\nabla p - p\nabla h) \cdot d\mathbf{n}.$$

To verify this, simply use the divergence theorem and product rule on the right-hand side of the equation. Since  $\Delta h = 0$  and  $p = 0$  on  $\partial R$ , Green's identity reduces to

$$\iint_R h\Delta p dA = \oint_{\partial R} h\nabla p \cdot d\mathbf{n} = \frac{d}{dt} \iint_R h dA.$$

Finally, by definition of the Green function, the area integral on the left is simply  $h(0)$ . □

5. Learn about Stokes' theorem for  $\mathbb{C}$ :

$$\iint_R \frac{\partial f}{\partial \bar{z}} d\bar{z} \wedge dz = \oint_{\partial R} f dz$$

and use it to prove the Cauchy integral theorem. Discuss.

*Proof.* Literally an entire book can be written for this. We'll just stick to the basics; from Stokes' theorem for differential forms,

$$\oint_{\partial R} f dz = \iint_R d(f dz) = \iint_R \frac{\partial f}{\partial \bar{z}} d\bar{z} \wedge dz + \underbrace{\frac{\partial f}{\partial z} dz \wedge dz}_{=0}$$

The requirements on  $R$  and  $f$  are that  $f$  must be continuously differentiable and  $R$  must have a boundary with continuously differentiable parametrization. A basic result in complex analysis is that complex differentiable functions  $f$  must satisfy  $\frac{\partial f}{\partial x} = -i\frac{\partial f}{\partial y}$ . In terms of  $\bar{z}$ , this is equivalent to

$$\frac{\partial f}{\partial \bar{z}} = \frac{1}{2} \left( \frac{\partial f}{\partial x} + i \frac{\partial f}{\partial y} \right) = 0.$$

So if  $f$  is complex differentiable (nontrivially, this is equivalent to 'analytic' when dealing with  $\mathbb{C}$ ), then

$$0 = \iint_R \frac{\partial f}{\partial \bar{z}} d\bar{z} \wedge dz = \oint_{\partial R} f dz.$$

This is the Cauchy integral theorem. Cauchy's theorem can actually be proved without requiring that  $f$  be continuously differentiable, but that requires an entirely different approach.  $\square$